Muon decay to one-loop order

in the LR-symmetric model

1. Introduction to the LRM
2. Basic assumptions for calculations
3. Renormalization of the model
4. Results and discussion
5. Conclusions and Outlook

M. Czakon, J. Gluza, M. Zrałek, Nucl. Phys. B573, 57 (2000),

M. Czakon, J. Gluza, F. Jegerlehner, M. Zrałek, Eur. Phys. J. C13, 275 (2000).

M. Czakon, J. Gluza, J. Hejczyk, hep-ph/0205303, to appear in Nucl. Phys. B.

1. Introduction to the LRM

The model has been constructed in 1973-1974,

Pati, Salam, Senjanovic, Mohapatra

$SU(2)_L \otimes SU(2)_R \otimes U(1)_{B-L}$ gauge group

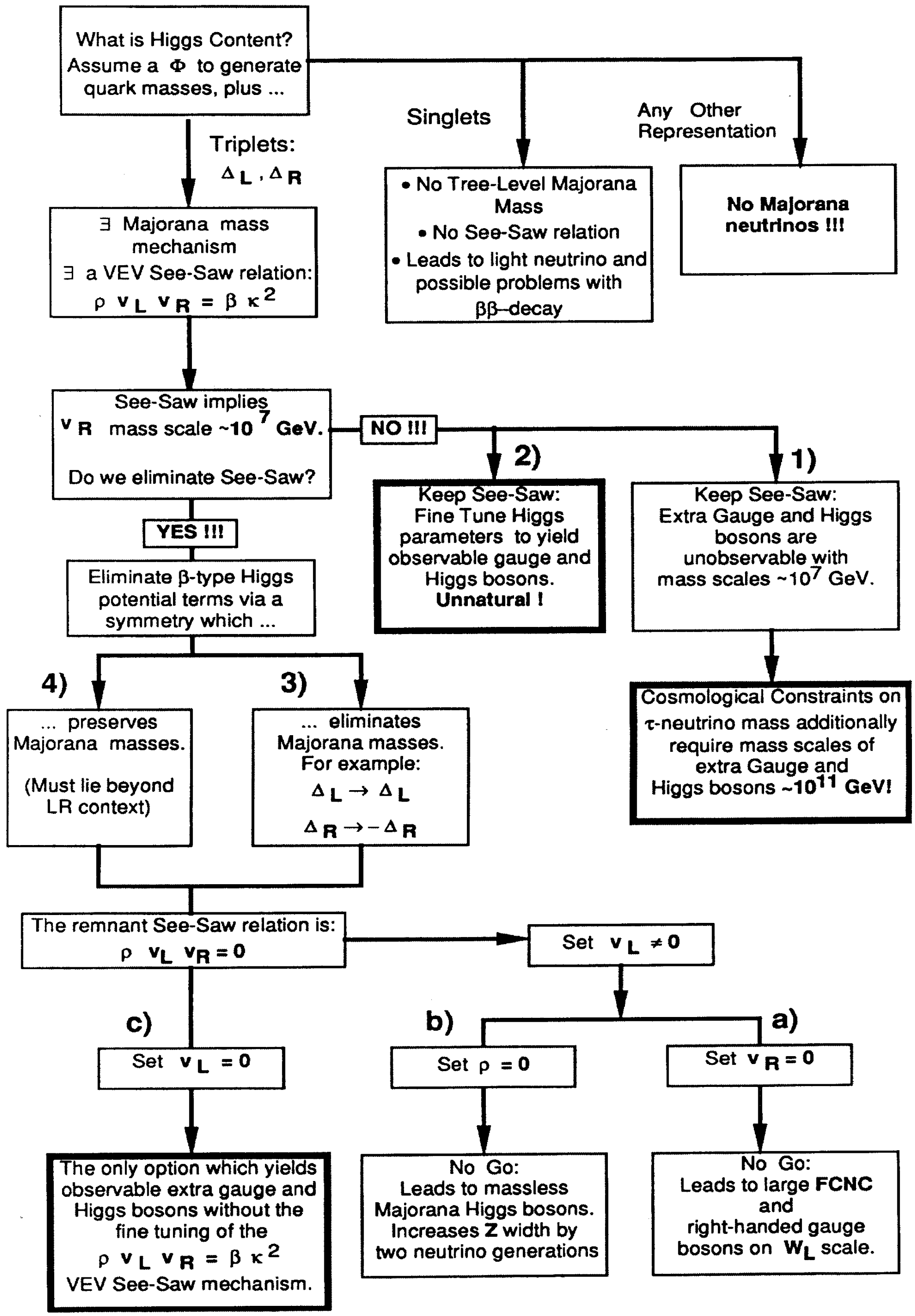
(i) restore lepton-quark symmetry to the weak interactions:

$$\begin{pmatrix} \nu_L \\ e_L \end{pmatrix}, \quad \begin{pmatrix} \nu_R \\ e_R \end{pmatrix}, \quad \begin{pmatrix} u_L \\ d_L \end{pmatrix}, \quad \begin{pmatrix} u_R \\ d_R \end{pmatrix}$$

(ii) Interpretation of the hypercharge as the difference of the baryon and lepton numbers:

$$Q = T_{3L} + T_{3R} + \frac{B-L}{2}$$

$$\begin{array}{lcl} W_L^{\pm}, W_L^0 & & W_1^{\pm}, W_2^{\pm} \\ W_R^{\pm}, W_R^0 & \rightarrow [SSB?] & Z_1, Z_2 \\ B^0 & & \gamma \end{array}$$


What is Higgs Content? Assume a Φ to generate quark masses, plus ...
Triplets: ΔL , ΔR
Singlets
Any Other Representation
∃ Majorana mass mechanism ∃ a VEV See-Saw relation: ρ vL vR = β κ2
• No Tree-Level Majorana Mass • No See-Saw relation • Leads to light neutrino and possible problems with ββ–decay
No Majorana neutrinos !!!
See-Saw implies vR mass scale ~10^7 GeV. Do we eliminate See-Saw?
NO !!!
2)
1)
Keep See-Saw: Fine Tune Higgs parameters to yield observable gauge and Higgs bosons. Unnatural !
Keep See-Saw: Extra Gauge and Higgs bosons are unobservable with mass scales ~10^7 GeV.
YES !!!
Eliminate β-type Higgs potential terms via a symmetry which ...
Cosmological Constraints on τ-neutrino mass additionally require mass scales of extra Gauge and Higgs bosons ~10^11 GeV!
4)
3)
... preserves Majorana masses. (Must lie beyond LR context)
... eliminates Majorana masses. For example: ΔL → ΔL ΔR → −ΔR
The remnant See-Saw relation is: ρ vL vR = 0
Set vL ≠ 0
c)
b)
a)
Set vL = 0
Set ρ = 0
Set vR = 0
The only option which yields observable extra gauge and Higgs bosons without the fine tuning of the ρ vL vR = β κ2 VEV See-Saw mechanism.
No Go: Leads to massless Majorana Higgs bosons. Increases Z width by two neutrino generations
No Go: Leads to large FCNC and right-handed gauge bosons on WL scale.

Deshpande, Gunion, Kayser, Olness, 1991

$$
\begin{aligned}
&\mathcal{L}_{Higgs} = \\
&-\mu_1^2 Tr[\Phi^\dagger \Phi] - \mu_2^2 (Tr[\tilde{\Phi}\Phi^\dagger] + Tr[\tilde{\Phi}^\dagger \Phi]) \\
&-\mu_3^2 (Tr[\Delta_L \Delta_L^\dagger] + Tr[\Delta_R \Delta_R^\dagger]) \\
&+\lambda_1 Tr[\Phi\Phi^\dagger]^2 + \lambda_2 (Tr[\tilde{\Phi}\Phi^\dagger]^2 + Tr[\tilde{\Phi}^\dagger \Phi]^2) \\
&+\lambda_3 (Tr[\tilde{\Phi}\Phi^\dagger] Tr[\tilde{\Phi}^\dagger \Phi]) \\
&+\lambda_4 (Tr[\Phi\Phi^\dagger](Tr[\tilde{\Phi}\Phi^\dagger] + Tr[\tilde{\Phi}^\dagger \Phi])) \\
&+\rho_1 (Tr[\Delta_L \Delta_L^\dagger]^2 + Tr[\Delta_R \Delta_R^\dagger]^2) \\
&+\rho_2 (Tr[\Delta_L \Delta_L] Tr[\Delta_L^\dagger \Delta_L^\dagger] + Tr[\Delta_R \Delta_R] Tr[\Delta_R^\dagger \Delta_R^\dagger]) \\
&+\rho_3 (Tr[\Delta_L \Delta_L^\dagger] Tr[\Delta_R \Delta_R^\dagger]) \\
&+\rho_4 (Tr[\Delta_L \Delta_L] Tr[\Delta_R^\dagger \Delta_R^\dagger] + Tr[\Delta_R \Delta_R] Tr[\Delta_L^\dagger \Delta_L^\dagger]) \\
&+\alpha_1 (Tr[\Phi\Phi^\dagger](Tr[\Delta_L \Delta_L^\dagger] + Tr[\Delta_R \Delta_R^\dagger])) \\
&+\alpha_2 (Tr[\Phi\tilde{\Phi}^\dagger] Tr[\Delta_R \Delta_R^\dagger] + Tr[\tilde{\Phi}\Phi^\dagger] Tr[\Delta_L \Delta_L^\dagger])) \\
&+\alpha_2^* (Tr[\Phi^\dagger\tilde{\Phi}] Tr[\Delta_R \Delta_R^\dagger] + Tr[\tilde{\Phi}^\dagger\Phi] Tr[\Delta_L \Delta_L^\dagger])) \\
&+\alpha_3 (Tr[\Phi\Phi^\dagger \Delta_L \Delta_L^\dagger] + Tr[\Phi^\dagger \Phi \Delta_R \Delta_R^\dagger]) \\
&+\beta_1 (Tr[\Phi \Delta_R \Phi^\dagger \Delta_L^\dagger] + Tr[\Phi^\dagger \Delta_L \Phi \Delta_R^\dagger]) \\
&+\beta_2 (Tr[\tilde{\Phi} \Delta_R \Phi^\dagger \Delta_L^\dagger] + Tr[\tilde{\Phi}^\dagger \Delta_L \Phi \Delta_R^\dagger]) \\
&+\beta_3 (Tr[\Phi \Delta_R \tilde{\Phi}^\dagger \Delta_L^\dagger] + Tr[\Phi^\dagger \Delta_L \tilde{\Phi} \Delta_R^\dagger]),
\end{aligned}
$$

invariant under the symmetry $\Delta_L \leftrightarrow \Delta_R, \quad \Phi \leftrightarrow \Phi^\dagger, \quad \beta_i = 0$.

The minimal Higgs sector consists of two triplets and one bidoublet

$$\Delta_{L,R} = \begin{pmatrix} \delta^{+}_{L,R}/\sqrt{2} & \delta^{++}_{L,R} \\ \delta^{0}_{L,R} & -\delta^{+}_{L,R}/\sqrt{2} \end{pmatrix},$$

$$\Phi = \begin{pmatrix} \phi^{0}_{1} & \phi^{+}_{1} \\ \phi^{-}_{2} & \phi^{0}_{2} \end{pmatrix}.$$

with vacuum expectation values allowed for the neutral particles

$$\frac{v_L}{\sqrt{2}} = \langle \delta^0_L \rangle,$$

$$\text{new HE scale}: \frac{v_R}{\sqrt{2}} = \langle \delta^0_R \rangle,$$

$$\text{SM VEV scale}: \sqrt{\kappa_1^2 + \kappa_2^2}$$

$$\frac{\kappa_1}{\sqrt{2}} = \langle \phi^0_1 \rangle,$$

$$\frac{\kappa_2}{\sqrt{2}} = \langle \phi^0_2 \rangle.$$

- The result is 20 real scalar fields, of which 14 are physical (the rest are Goldstone bosons):
  - 4 neutral scalars: $H_0^0, H_1^0, H_2^0, H_4^0$,
    (the first can be considered to be the light Higgs of the SM at tree level),
  - 2 neutral pseudo-scalars: A_1^0, A_2^0,
  - 2 charged scalars: $H_1^\pm, H_2^\pm$,
  - 2 doubly-charged scalars: $H_1^{\pm\pm}, H_2^{\pm\pm}$.
- see-saw mechanism for the generation of light neutrino masses, with specific SB sectors. The neutrino mass matrix

$$M_\nu = \begin{pmatrix} M_L(v_L) & M_D(\kappa_{1,2}) \\ M_D^T & M_R(v_R) \end{pmatrix}$$

with $M_L \ll M_D \ll M_R$.

$$\begin{aligned} m_N &\sim M_R \\ m_{\rm light} &\sim M_D^2/M_R \end{aligned}$$

$M_D \sim \mathcal{O}(1)\ \mathrm{GeV} \to \mathrm{M_R} \sim 10^{15}\ \mathrm{GeV}$, if light neutrino masses of the order of 0.1 eV.

W_R (Right-Handed W Boson) MASS LIMITS

Assuming a light right-handed neutrino, except for BEALL 82, LANGACKER 89B, and COLANGELO 91. $g_R = g_L$ assumed. [Limits in the section MASS LIMITS for W' below are also valid for W_R if $m_{\nu_R} \ll m_{W_R}$.] Some limits assume manifest left-right symmetry, *i.e.*, the equality of left- and right Cabibbo-Kobayashi-Maskawa matrices. For a comprehensive review, see LANGACKER 89B. Limits on the W_L-W_R mixing angle ζ are found in the next section. Values in brackets are from cosmological and astrophysical considerations and assume a light right-handed neutrino.

VALUE (GeV)	*CL%*	*DOCUMENT ID*		*TECN*	*COMMENT*
> 715	90	[11] CZAKON	99	RVUE	Electroweak
> 310	90	[12] THOMAS	01	CNTR	β^+ decay
> 137	95	[13] ACKERSTAFF	99D	OPAL	τ decay
>1400	68	[14] BARENBOIM	98	RVUE	Electroweak, Z-Z' mixing
> 549	68	[15] BARENBOIM	97	RVUE	μ decay
> 220	95	[16] STAHL	97	RVUE	τ decay
> 220	90	[17] ALLET	96	CNTR	β^+ decay
> 281	90	[18] KUZNETSOV	95	CNTR	Polarized neutron decay
> 282	90	[19] KUZNETSOV	94B	CNTR	Polarized neutron decay
> 439	90	[20] BHATTACH...	93	RVUE	Z-Z' mixing
> 250	90	[21] SEVERIJNS	93	CNTR	β^+ decay
		[22] IMAZATO	92	CNTR	K^+ decay
> 475	90	[23] POLAK	92B	RVUE	μ decay
> 240	90	[24] AQUINO	91	RVUE	Neutron decay
> 496	90	[24] AQUINO	91	RVUE	Neutron and muon decay
> 700		[25] COLANGELO	91	THEO	$m_{K^0_L} - m_{K^0_S}$
> 477	90	[26] POLAK	91	RVUE	μ decay
[none 540–23000]		[27] BARBIERI	89B	ASTR	SN 1987A; light ν_R
> 300	90	[28] LANGACKER	89B	RVUE	General
> 160	90	[29] BALKE	88	CNTR	$\mu \to e\nu\overline{\nu}$
> 406	90	[30] JODIDIO	86	ELEC	Any ζ
> 482	90	[30] JODIDIO	86	ELEC	$\zeta = 0$
> 800		MOHAPATRA	86	RVUE	$SU(2)_L \times SU(2)_R \times U(1)$
> 400	95	[31] STOKER	85	ELEC	Any ζ
> 475	95	[31] STOKER	85	ELEC	$\zeta < 0.041$
		[32] BERGSMA	83	CHRM	$\nu_\mu e \to \mu \nu_e$
> 380	90	[33] CARR	83	ELEC	μ^+ decay
>1600		[34] BEALL	82	THEO	$m_{K^0_L} - m_{K^0_S}$
[> 4000]		STEIGMAN	79	COSM	Nucleosynthesis; light ν_R

3. Assumptions

Tree level diagrams to the muon decay:

- exchanged W_1,
- exchanged W_2,
- exchanged charged Higgs particles.

We assume:

- no mixing in the charged gauge sector: $\kappa_2 = 0, \quad \xi = 0$,
- mixings among light and heavy neutrino states are neglected

Effects of heavy neutrinos remain:

- $Z_1 N_i N_i,\ Z_2 N_i N_i$ (build up gauge-boson self energies);
- $\bar{e} N_4 H_2^-$: $m_{N_4} \frac{g}{2M_{W_2}} \frac{\kappa_1}{v_R} P_L$,
- $\bar{e} N_1 H_1^-$: $\sqrt{2}\frac{m_{N_4}}{v_R} P_R$

Nice feature: constraints on the heavy (right) sector come directly from 1-loop level!

$$
\begin{aligned}
-\mathcal{L} &= \bar{c}_{LL}\bar{e}\gamma_\alpha P_L \nu_e \bar{\nu}_\mu \gamma^\alpha P_L \mu + \bar{c}_{RR}\bar{e}\gamma_\alpha P_R \nu_e \bar{\nu}_\mu \gamma^\alpha P_R \mu \\
&+ \bar{c}_{LR}\bar{e}\gamma_\alpha P_L \nu_e \bar{\nu}_\mu \gamma^\alpha P_R \mu + \bar{c}_{RL}\bar{e}\gamma_\alpha P_R \nu_e \bar{\nu}_\mu \gamma^\alpha P_L \mu.
\end{aligned}
$$

where:

$$
\begin{aligned}
\bar{c}_{LL} &= \frac{g^2}{2M_{W_1}^2}(\cos^2\xi + \beta\sin^2\xi), \\
\bar{c}_{RR} &= \frac{g^2}{2M_{W_1}^2}(\sin^2\xi + \beta\cos^2\xi), \\
\bar{c}_{RL} = \bar{c}_{LR} &= \frac{g^2}{2M_{W_1}^2}(-1+\beta)\sin\xi\cos\xi.
\end{aligned}
$$

$\beta = \frac{M_{W_1}^2}{M_{W_2}^2}$,

To have neutrino mixings properly included:

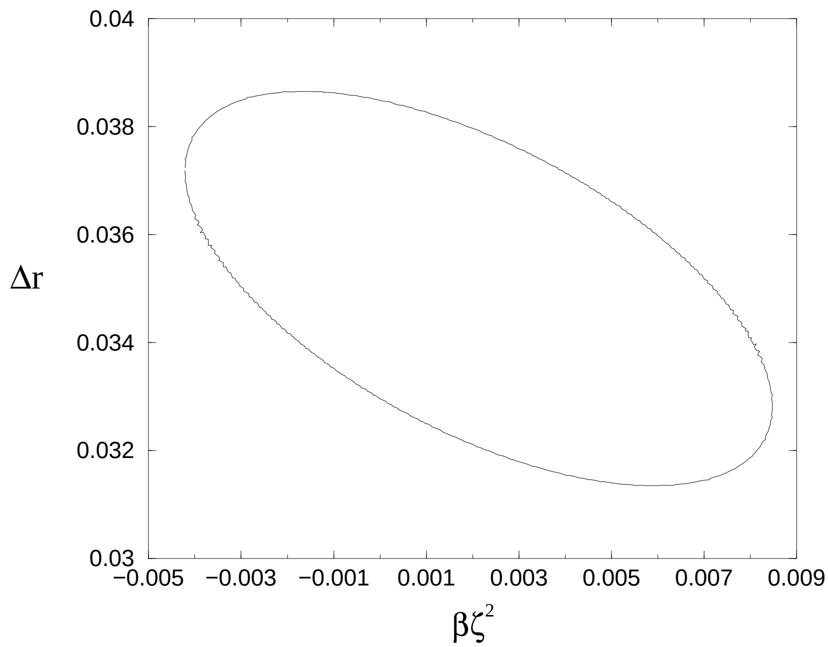
$$
U = \begin{pmatrix} K_L^T \\ K_R^\dagger \end{pmatrix} = \begin{pmatrix} O(1) & O(1/M_N) \\ O(1/M_N) & O(1) \end{pmatrix}.
$$

$$
\begin{aligned}
c_{LL} &= \bar{c}_{LL}(K_L^\dagger)_{ea}(K_L)_{b\mu}, \\
c_{RR} &= \bar{c}_{RR}(K_R^\dagger)_{ea}(K_R)_{b\mu}, \\
c_{LR} &= \bar{c}_{LR}(K_L^\dagger)_{ea}(K_R)_{b\mu}, \\
c_{RL} &= \bar{c}_{RL}(K_R^\dagger)_{ea}(K_L)_{b\mu}.
\end{aligned}
$$

then: $c_{LL} \gg c_{RR}, c_{LR}, c_{RL}$

and

$$\frac{G_F}{\sqrt{2}} \simeq |c_{LL}|^2 \quad \simeq \quad \frac{\pi\alpha}{2s_W^2 M_W^2 (1-\Delta r)}(1+\beta\zeta^2).$$



with the very worst bound $\zeta < .1$, it gives $\beta < 1 \quad \rightarrow M_{W_2} \geq M_{W_1}$

without heavy neutrino mixing:

$$\begin{aligned} \frac{G_F}{\sqrt{2}} &= |c_{LL}|^2 + |c_{LR}|^2 + |c_{RL}|^2 + |c_{RR}|^2 \\ &= \frac{\pi\alpha}{2s_W^2 M_W^2, (1-\Delta_R)}(1+\beta^2) \end{aligned}$$

and $M_{W_2} \geq 500 \text{ GeV}$

3. Gauge Sector Renormalization

- The SSB leads to the following gauge-boson mass matrices

$$\frac{g^2}{4}\begin{pmatrix} \kappa_+^2 & -2\kappa_1\kappa_2 \\ -2\kappa_1\kappa_2 & \kappa_+^2 + 2v_R^2 \end{pmatrix},$$

$$\frac{1}{2}\begin{pmatrix} \frac{1}{2}g^2\kappa_+^2 & -\frac{1}{2}g^2\kappa_+^2 & 0 \\ -\frac{1}{2}g^2\kappa_+^2 & \frac{1}{2}g^2(\kappa_+^2 + 4v_R^2) & -2gg'v_R^2 \\ 0 & -2gg'v_R^2 & 2g'^2v_R^2 \end{pmatrix}$$

- The diagonalization introduces various mixing angles. It is possible to write the respective matrices as

$$\begin{pmatrix} W_L^\pm \\ W_R^\pm \end{pmatrix} = \begin{pmatrix} \cos\xi & \sin\xi \\ -\sin\xi & \cos\xi \end{pmatrix}\begin{pmatrix} W_1^\pm \\ W_2^\pm \end{pmatrix},$$

$$\begin{pmatrix} W_L^3 \\ W_R^3 \\ B \end{pmatrix} = \begin{pmatrix} c_W c & c_W s & s_W \\ -s_W s_M c - c_M s & -s_W s_M s + c_M c & c_W s_M \\ -s_W c_M c + s_M s & -s_W c_M s - s_M c & c_W c_M \end{pmatrix}\begin{pmatrix} Z_1 \\ Z_2 \\ A \end{pmatrix},$$

where:

$$c_W \equiv \cos\theta_W, \quad s_W \equiv \sin\theta_W, \quad s \equiv \sin\phi,$$

$$c_M \equiv \frac{\sqrt{\cos 2\theta_W}}{\cos\theta_W}, \quad c \equiv \cos\phi. \quad s_M \equiv \tan\theta_W$$

- $$M^2_{W_{1,2}} = \frac{g^2}{4}\left(\kappa_+^2 + v_R^2 \mp \sqrt{v_R^4 + 4\kappa_1^2\kappa_2^2}\right),$$
$$M^2_{Z_{1,2}} = \frac{1}{4}\left(\left((g^2\kappa_+^2 + 2v_R^2(g^2+{g'}^2)\right)\right.$$
$$\left.\mp\sqrt{(g^2\kappa_+^2 + 2v_R^2(g^2+{g'}^2))^2 - 4g^2(g^2+2{g'}^2)\kappa_+^2 v_R^2}\right)$$

  and

  $$\tan 2\xi = -\frac{2\kappa_1\kappa_2}{v_R^2}, \qquad \sin 2\phi = -\frac{g^2\kappa_+^2\sqrt{\cos 2\theta_W}}{2\cos^2\theta_W\,(M^2_{Z_2} - M^2_{Z_1})}.$$
$$g = \frac{e}{\sin\theta_W}, \qquad g' = \frac{e}{\sqrt{\cos 2\theta_W}}.$$

- We exploit the one-to-one correspondence of the parameters:

$$\begin{array}{ccc} g & & e, \\ g' & & M_{W_1}, \\ \kappa_1 & \rightarrow & M_{W_2}, \\ \kappa_2 & & M_{Z_1}, \\ v_R & & M_{Z_2}. \end{array}$$

  Adopt an on-shell scheme in which we renormalize the self-energies of the gauge-bosons by requiring that they vanish at the physical mass squared. The electric charge is renormalized as in the SM.

- The renormalization parameter that is of particular interest is the Weinberg angle.
  By expanding to first order, we end up with

$$\delta s_W^2 = 2c_W^2 \frac{(\delta M_{Z_2}^2 + \delta M_{Z_1}^2) - (\delta M_{W_2}^2 + \delta M_{W_1}^2)}{(M_{Z_2}^2 + M_{Z_1}^2) - (M_{W_2}^2 + M_{W_1}^2)}$$

$$+ \frac{1}{2} \frac{(M_{W_2}^2 + M_{W_1}^2)(\delta M_{Z_2}^2 + \delta M_{Z_1}^2) + (M_{Z_2}^2 + M_{Z_1}^2)(\delta M_{W_2}^2 + \delta M_{W_1}^2)}{\Big((M_{Z_2}^2 + M_{Z_1}^2) - (M_{W_2}^2 + M_{W_1}^2)\Big)^2}$$

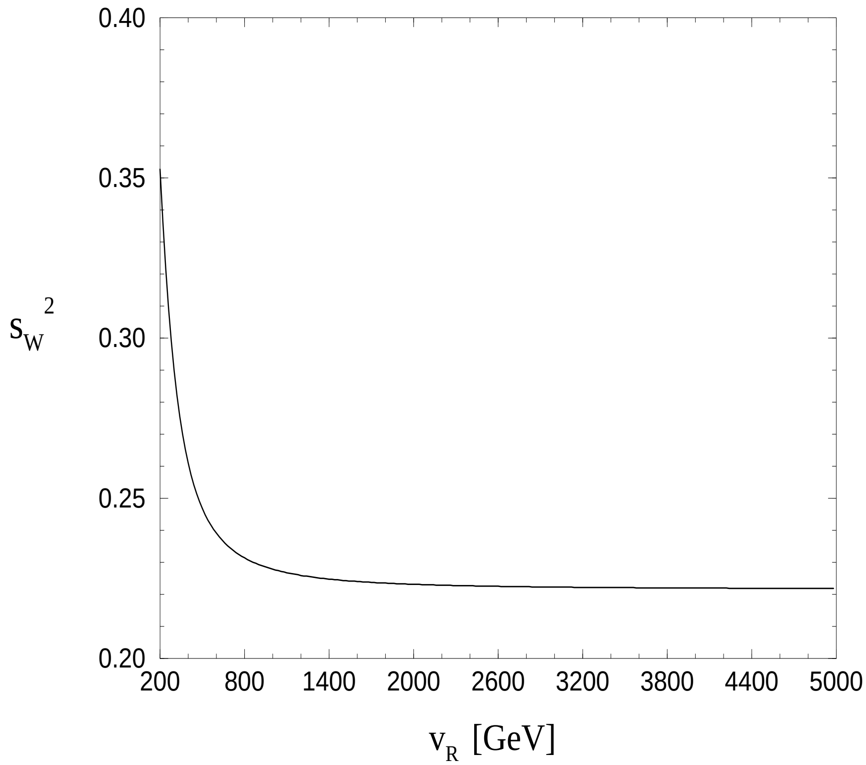
$$- \frac{1}{2} \frac{(2M_{Z_1}^2 + M_{Z_2}^2)\delta M_{Z_1}^2 + (2M_{Z_2}^2 + M_{Z_1}^2)\delta M_{Z_2}^2}{\Big((M_{Z_2}^2 + M_{Z_1}^2) - (M_{W_2}^2 + M_{W_1}^2)\Big)^2} .$$

  The denominator is expressible through the heavy breaking scale only

$$\Big((M_{Z_2}^2 + M_{Z_1}^2) - (M_{W_2}^2 + M_{W_1}^2)\Big) = \frac{g^2}{2c_M^2 c_W^2} v_R^2 .$$

  In the SM:

$$\delta s_W^2 = c_W^2 \left(\frac{\delta M_Z^2}{M_Z^2} - \frac{\delta M_W^2}{M_W^2} \right)$$



in our model s_W^2 is expressable by M_{W_1},M_{Z_1}, and a heavy scale v_R,

SM:

$$s_W^2 = 1 - \frac{M_{W_1}^2}{M_{Z_1}^2}$$

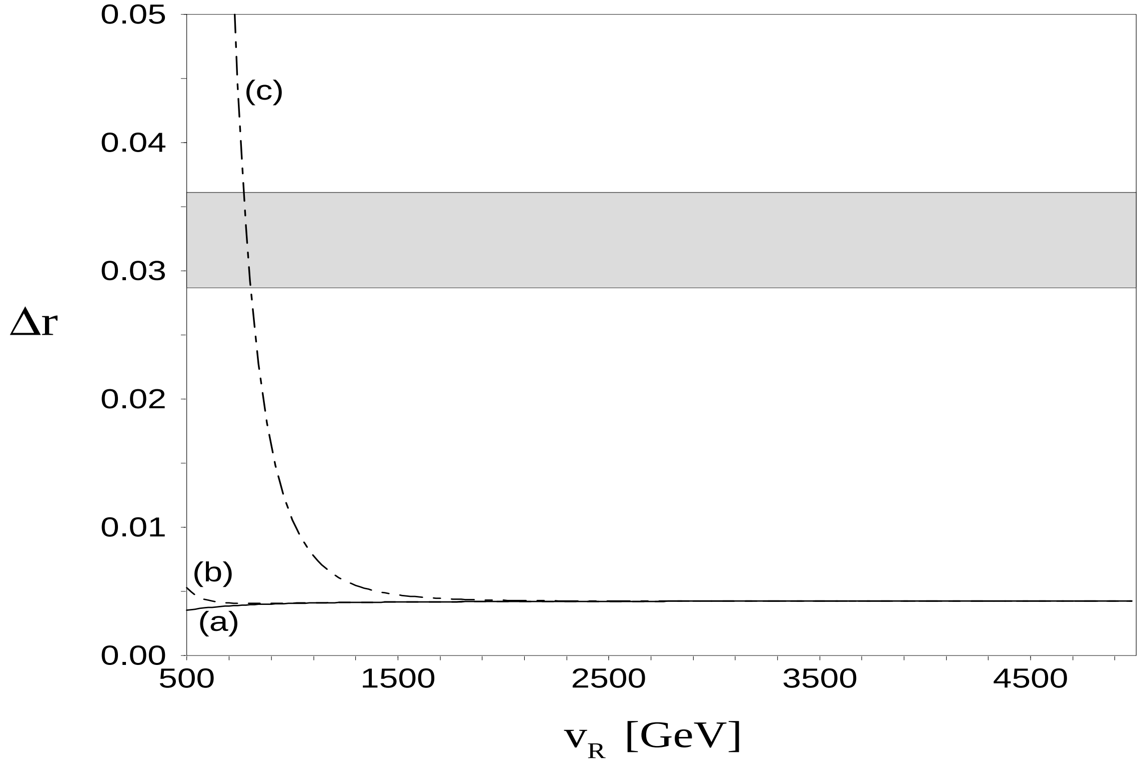
We include the change of s_W^2 from the SM to the LRM in Δr,

$$\frac{e^2}{(8M_W^2 s_W^2)_{SM}}(1+\Delta r) = \frac{e^2}{(8M_{W_1}^2 s_W^2)_{LR}}(1+\Delta r_{LR})$$

then:

$$\begin{aligned}
\Delta r &= \frac{(s_W^2)_{SM}}{(s_W^2)_{LRSM}}\left(\frac{-\Pi_W^T(0)-\delta M_W^2}{M_W^2}+2\frac{\delta e}{e}\right. \\
&- \left.\frac{\delta s_W^2}{s_W^2}+\delta_V+\delta_B\right) \\
&- \frac{(s_W^2)_{LRSM}-(s_W^2)_{SM}}{(s_W^2)_{LRSM}},
\end{aligned}$$

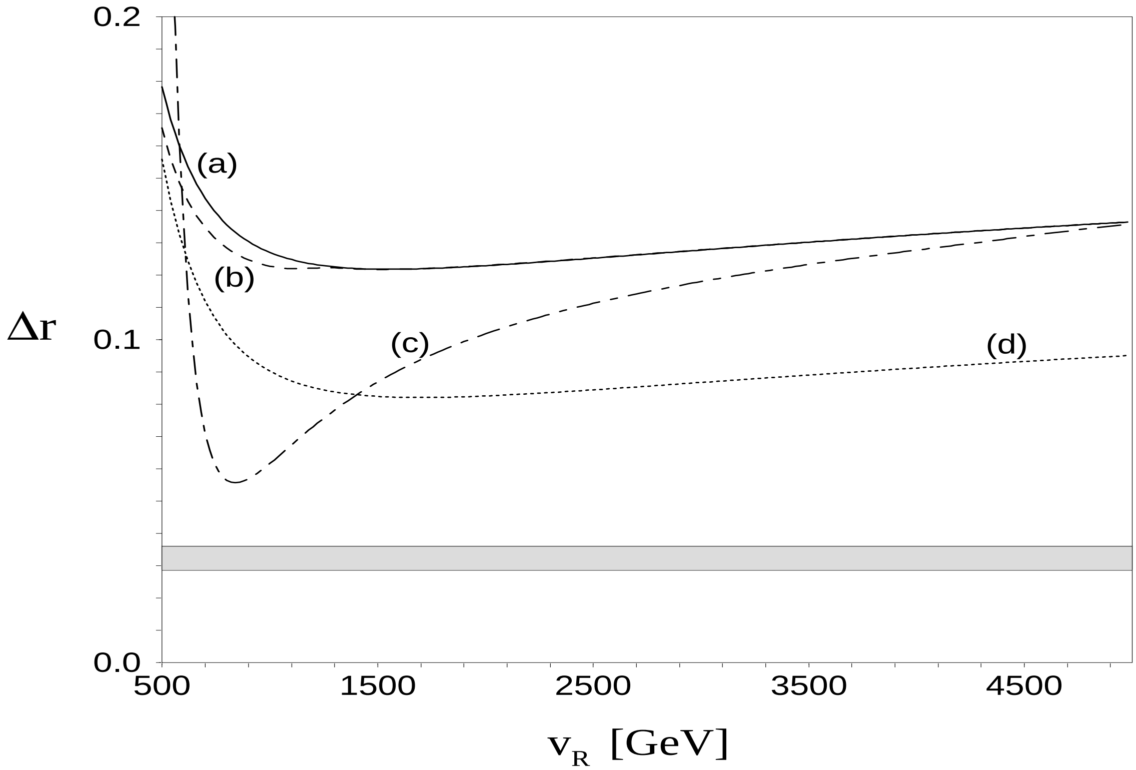
4. Results and Discussion



Contribution of box diagrams to Δr. With increasing v_R heavy particles decouple and the lines aim at the SM contribution. The (a) line is for (three heavy neutrinos) $m_N = 100$ GeV; (b) is for $m_N = 500$ GeV; (c) is for $m_N = 2$ TeV. The gray area shows the experimentally allowed values of Δr.

$$m_N = \sqrt{2} h_M v_R \leq \sqrt{2} v_R$$

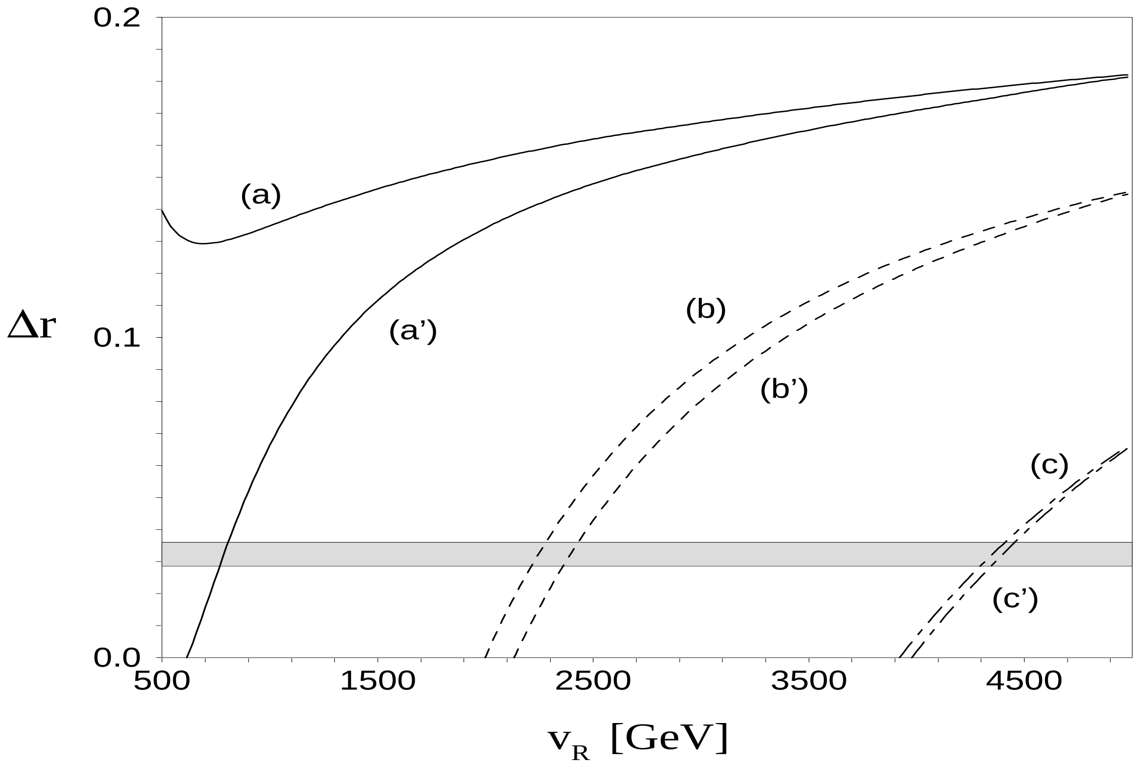
"Natural" scalar potential:



The (a) line is for (three heavy neutrinos) $m_N = 100$ GeV; (b) is for $m_N = 500$ GeV; (c) is for $m_N = 2$ TeV.

$$\begin{aligned} M_{H_a} &\equiv M_{H_1^0} = M_{H_3^0} = M_{A_1^0} = M_{A_2^0} \\ &= M_{H_1^+} = M_{H_2^+} = M_{\delta_L^{++}} = v_R/\sqrt{2}, \\ M_{H_b} &\equiv M_{H_2^0} = M_{\delta_R^{++}} = \sqrt{2} v_R, \\ M_{H_0^0} &= \sqrt{2}\kappa_1. \end{aligned}$$

All Higgs particle masses equal (apart from H_0^0):



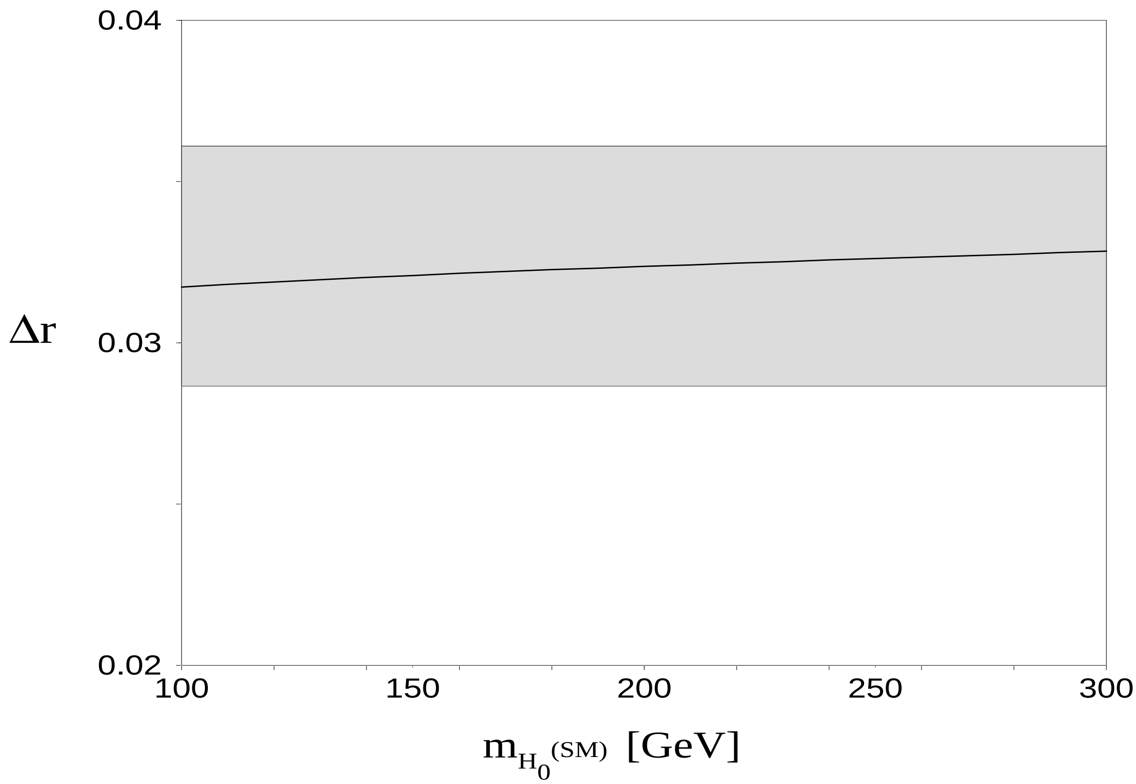
Sets with and without primes show results for three heavy neutrino masses

with $m_N = 100$ GeV and $m_N = 2$ TeV respectively.

for $M_H \simeq 2v_R$ (and large v_R), Δr can be accomodated,

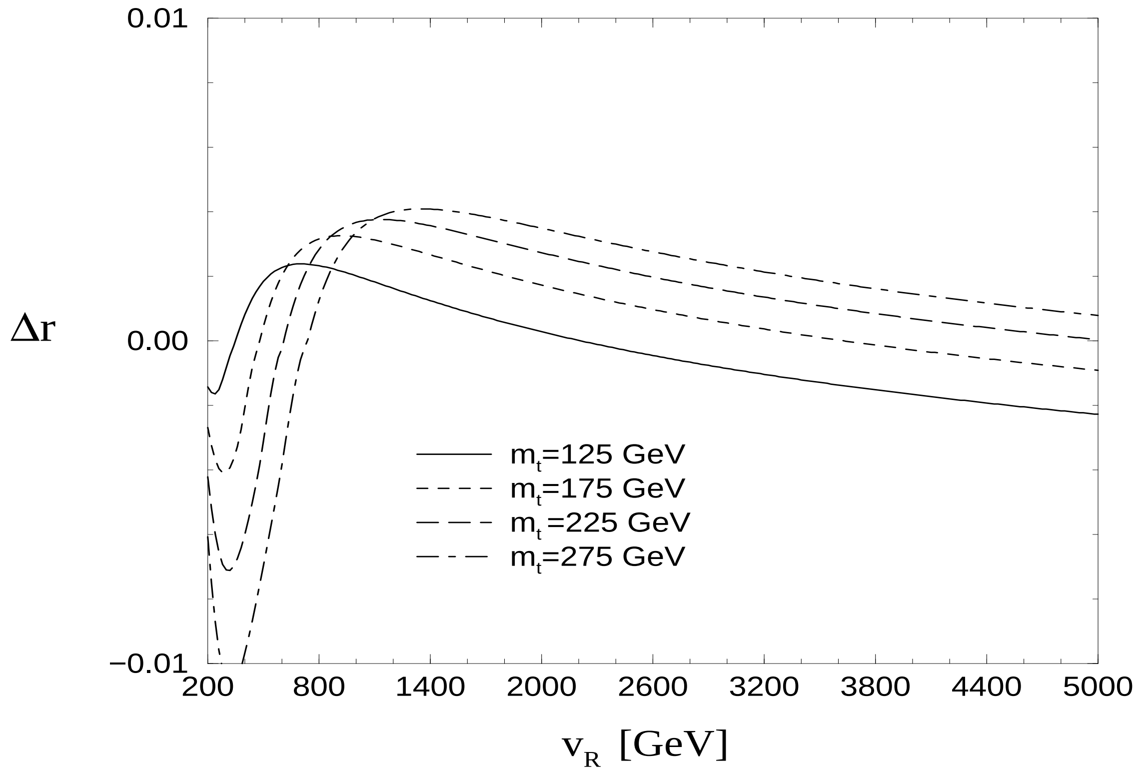
remark: fine tuned masses

Variaton of $m(H_0^0)$ is negligible:



Δr as function of the lightest Higgs scalar mass $M_{H_0^0}$. The gray area shows the experimentally allowed values of Δr and the heavy particle spectrum is chosen to fi t approximately to this region, namely, $v_R = 2390$ GeV, $m_N = 2$ TeV, $m_H = 5$ TeV.

m_t is effectively light and is not predictable by Δr in LRM:



The contribution of the third quark family to Δr as function of v_R for different top quark masses.

Leading terms:

$$(\Delta\rho)_{SM} \simeq \frac{m_t^2}{M_W^2}$$

$$(\Delta\rho)_{LRM} \simeq \frac{m_t^2}{M_{W_2}^2 - M_{W_1}^2}$$

6. Conclusions and Outlook

LRM can be fitted to muon decay width, however:

- heavy neutrino masses and Higgs masses must be very tuned to get proper Δr:

  $\rightarrow$ low energy data ? ...

- corrections coming from SM particles within LR model do not constitute a structure equivalent to their SM structure, e.g. top quark and $\Delta\rho$, the lightest Higgs;

- consequences:

  Czakon, Gluza, Jegerlehner, Zrałek, EPJ, C13, 275 (2000),

  also Lynn & Nardi, NPB381, 1992, ...

- LRM gives a possibility to study many effects: RH currents, $W_2, Z_2, \delta^{++}, \delta^{+}, N$, further questions in 1-loop level:

$$\begin{array}{ccc} g & & e, \\ g' & & M_{W_1}, \\ \kappa_1 & \rightarrow & M_{W_2}, \\ \kappa_2 & & M_{Z_1}, \\ v_R & & M_{Z_2}. \end{array}$$

$g_L \neq g_R, v_L \neq 0$, non-finite QED with $\xi \neq 0$,...